

Muon Monte Carlo: a new high-precision tool for muon propagation through matter

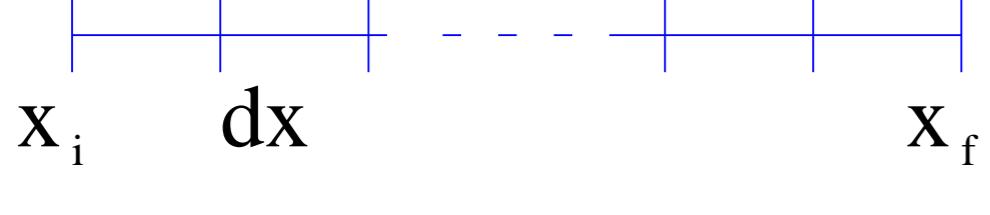
Dmitry Chirkin^{1,2}, Wolfgang Rhode²

¹University of California at Berkeley, USA

²Bergische Universität Gesamthochschule Wuppertal, Germany

Derivation of tracking formulae

Continuous part of the energy losses: $-\frac{dE}{dx} = f(E)$



The stochastic part: $dP(E(x_i)) = \sigma(E(x_i))dx$. Probability to suffer a catastrophic loss on dx at x_f is

$$\begin{aligned} & (1 - dP(E(x_i))) \cdot \dots \cdot (1 - dP(E(x_f))) \cdot dP(E(x_f)) \\ &= \exp(-dP(E(x_i))) \cdot \dots \cdot \exp(-dP(E(x_f))) \cdot dP(E(x_f)) \\ &= \exp\left(-\int_{E_i}^{E_f} dP(E(x))\right) \cdot dP(E(x_f)) \\ &= d_f \left(-\exp\left(-\int_{E_i}^{E_f} \frac{\sigma(E)}{-f(E)} \cdot dE\right)\right) = d(-\xi), \quad \xi \in (0; 1] \end{aligned}$$

The above equation can be solved for E_f :

$$\int_{E_i}^{E_f} \frac{\sigma(E)}{-f(E)} \cdot dE = -\log(\xi) \quad (\text{energy integral}),$$

and then the corresponding displacement is found:

$$x_f = x_i - \int_{E_i}^{E_f} \frac{dE}{f(E)} \quad (\text{tracking integral}).$$

Continuous randomization

choose dx so small that $p_0 = \int_{e_0}^{e_{cut}} p(e; E) de \cdot dx \ll 1$

$$\begin{aligned} < e > &= \int_{e_0}^{e_{cut}} e \cdot p(e; E) de \cdot dx \quad < e^2 > = \int_{e_0}^{e_{cut}} e^2 \cdot p(e; E) de \cdot dx \\ &< (\Delta e)^2 > = < e^2 > - < e >^2 \end{aligned}$$

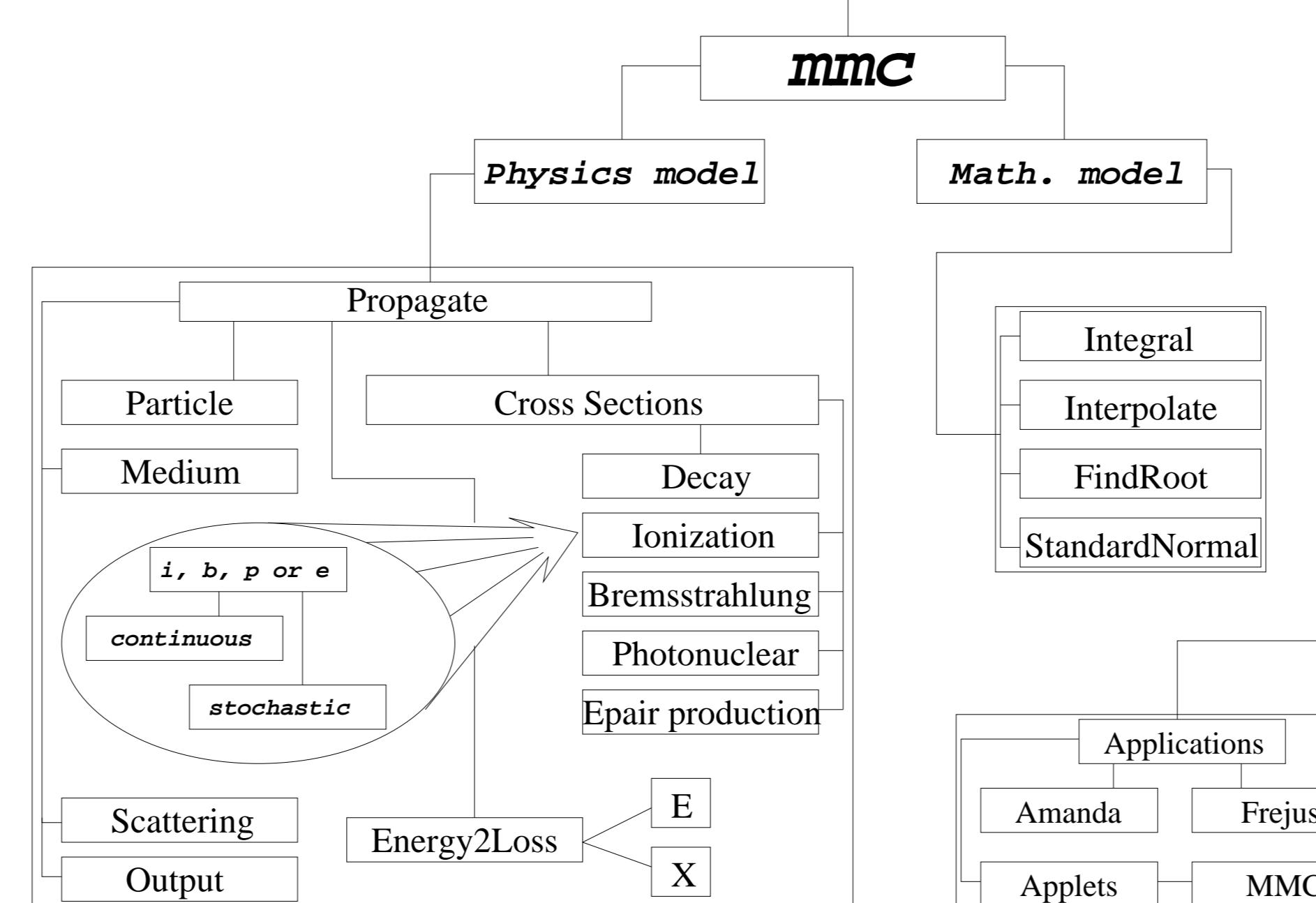
For small e_{cut} , $p(e; E) \simeq p(e; E_i) \rightarrow$ distributions $p(e; E)$ are the same \rightarrow central limit theorem can be applied

$$\begin{aligned} & < (\Delta(\Delta E))^2 > = \sum_n \left(< e_n^2 > - < e_n >^2 \right) = \\ & \sum_n \left[\left(\int_{e_0}^{e_{cut}} e_n^2 \cdot p(e_n; E_i) de_n \right) dx_n - \left(\int_{e_0}^{e_{cut}} e_n \cdot p(e_n; E_i) de_n \right)^2 dx_n^2 \right] \\ & \simeq \int_{x_i}^{x_f} dx \cdot \left[\left(\int_{e_0}^{e_{cut}} e^2 \cdot p(e; E(x)) de \right) - \left(\int_{e_0}^{e_{cut}} e \cdot p(e; E(x)) de \right)^2 dx \right] \end{aligned}$$

Here E_i was replaced with average expectation value of energy at $x, E(x)$. As $dx \rightarrow 0$, the second term disappears. The lower limit of the integral over e can be replaced with zero, since all of the cross sections diverge slower than $1/e^3$. Then,

$$< (\Delta(\Delta E))^2 > \simeq \int_{x_i}^{x_f} \frac{dE}{-f(E)} \cdot \left(\int_0^{e_{cut}} e^2 \cdot p(e; E) de \right)$$

MMC structure

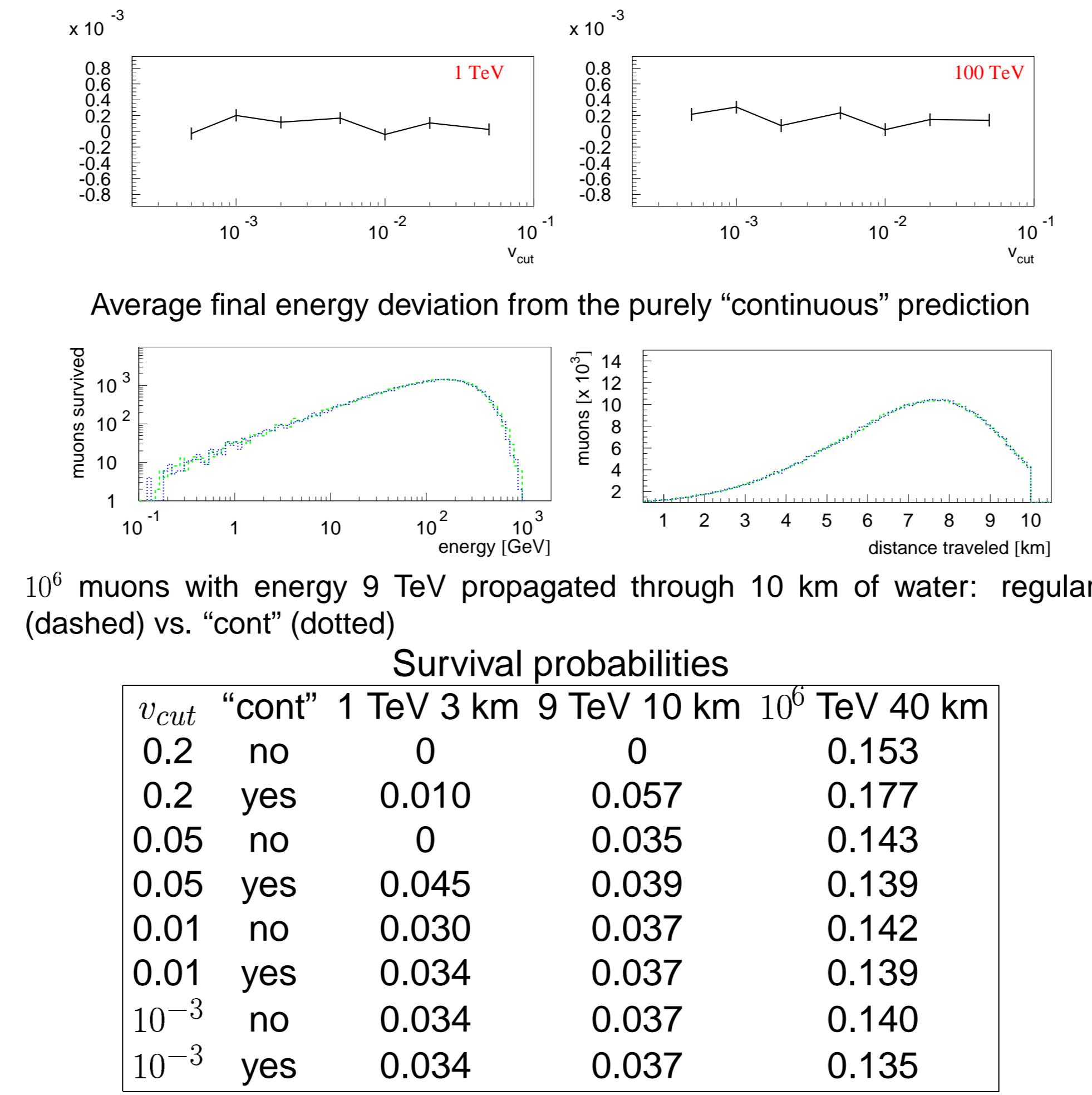


Results

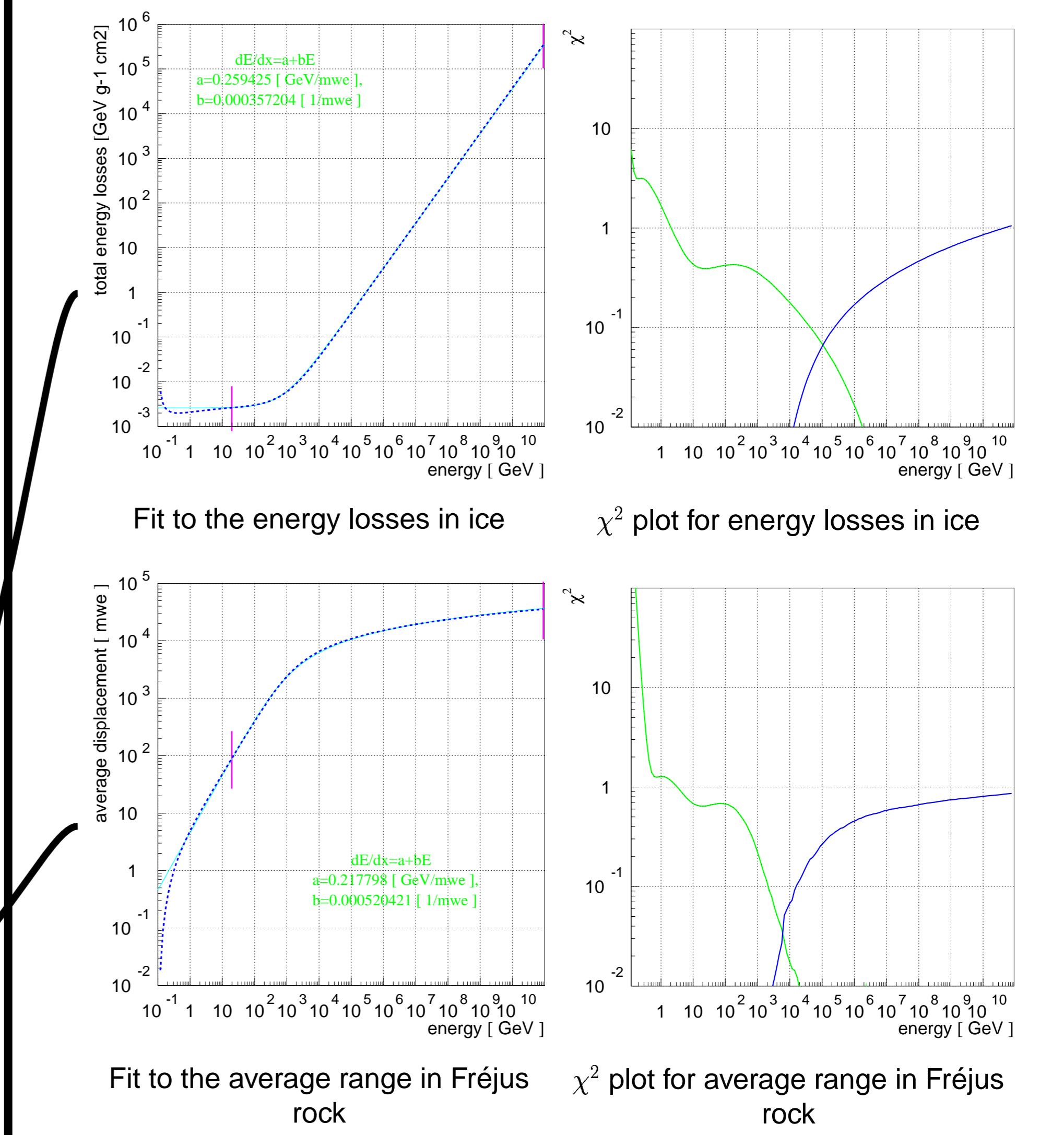
medium	a, GeV	b, mwe	av. dev.	max. dev.
ice	0.259	0.357	3.7%	6.6%
fr. rock	0.231	0.429	3.0%	5.1%

medium	a, GeV	b, mwe	av. dev.
ice	0.268	0.470	3.0%
fréjus rock	0.218	0.520	2.8%

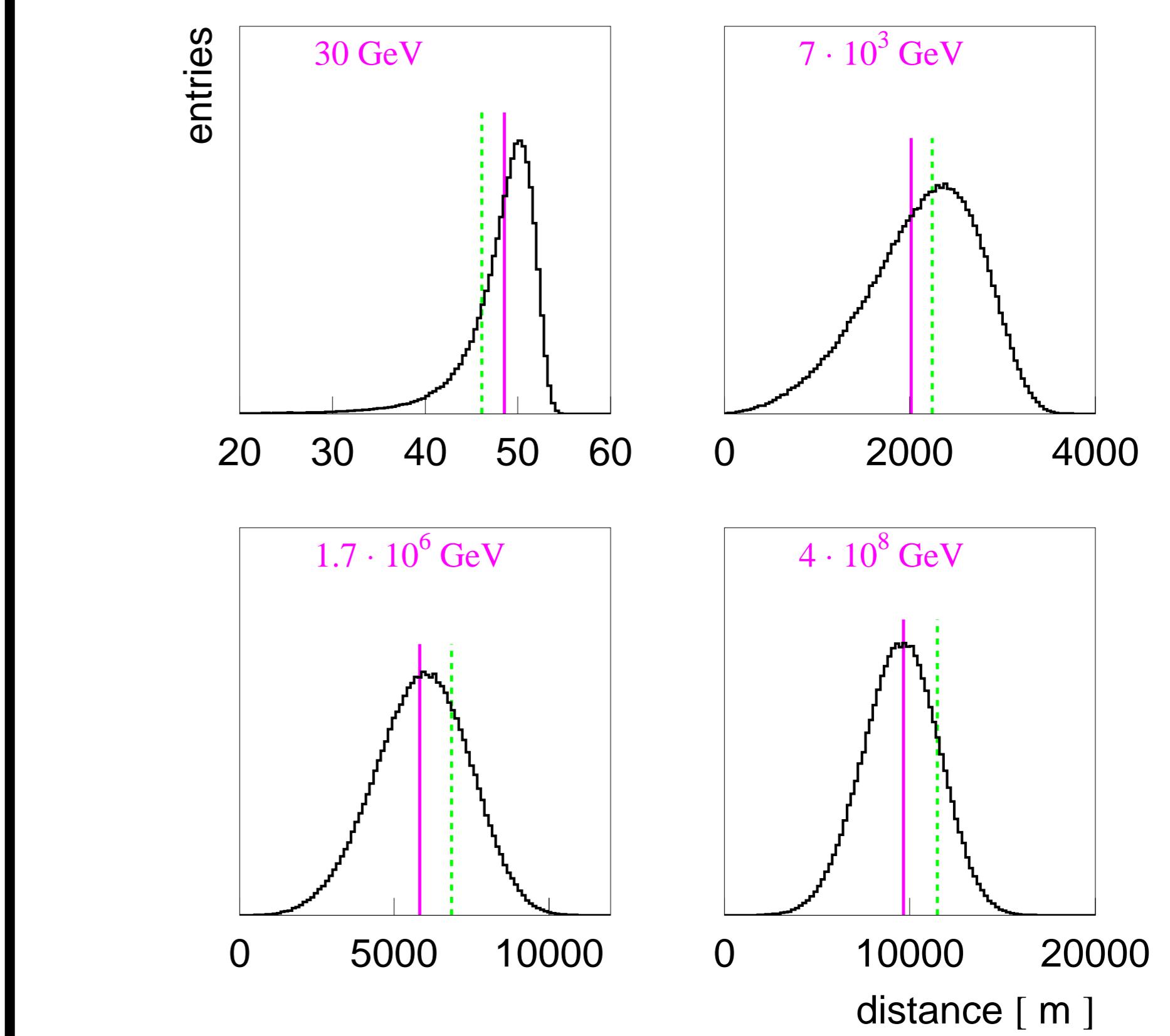
Algorithm errors



Fits to total energy loss and average range

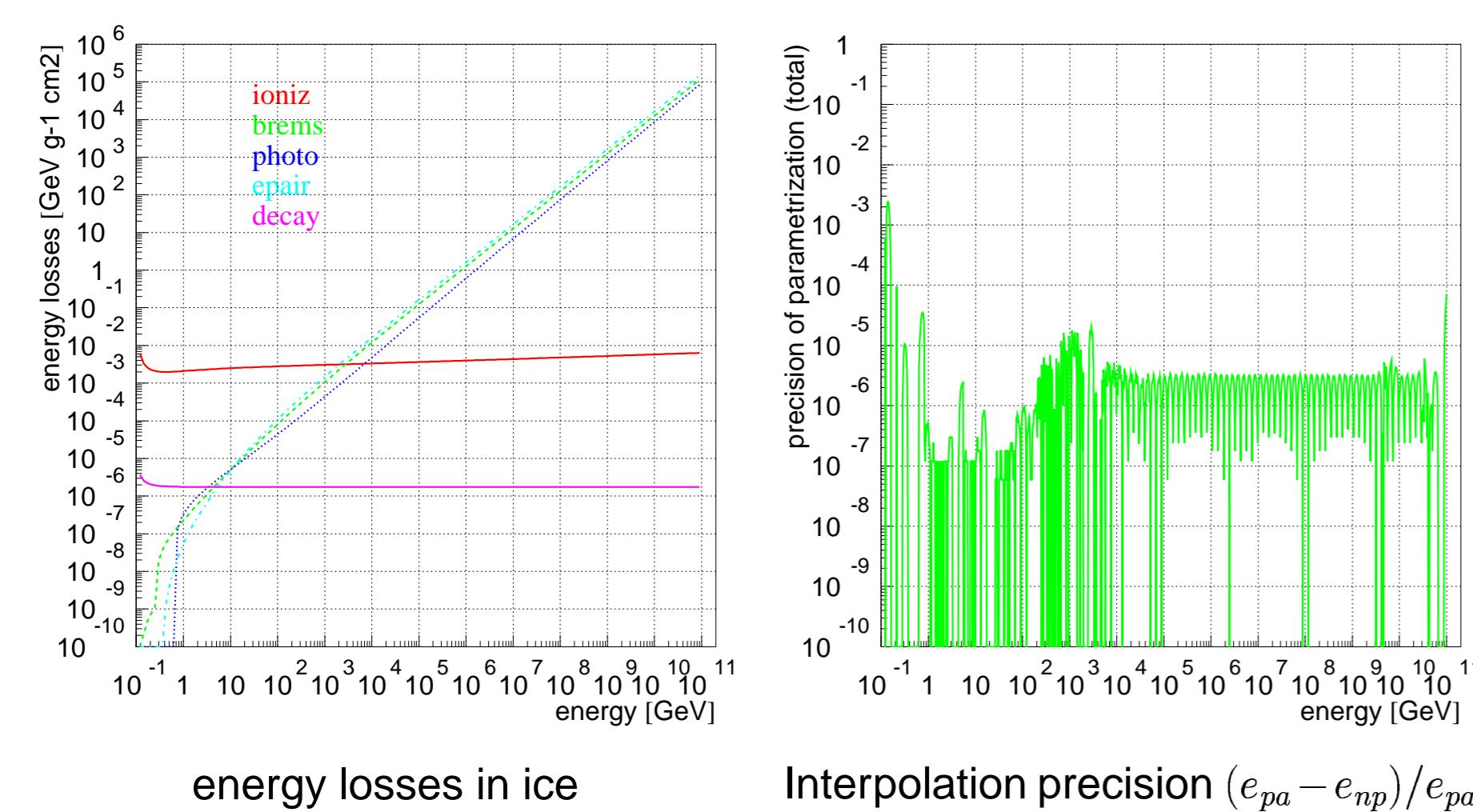


Range distributions

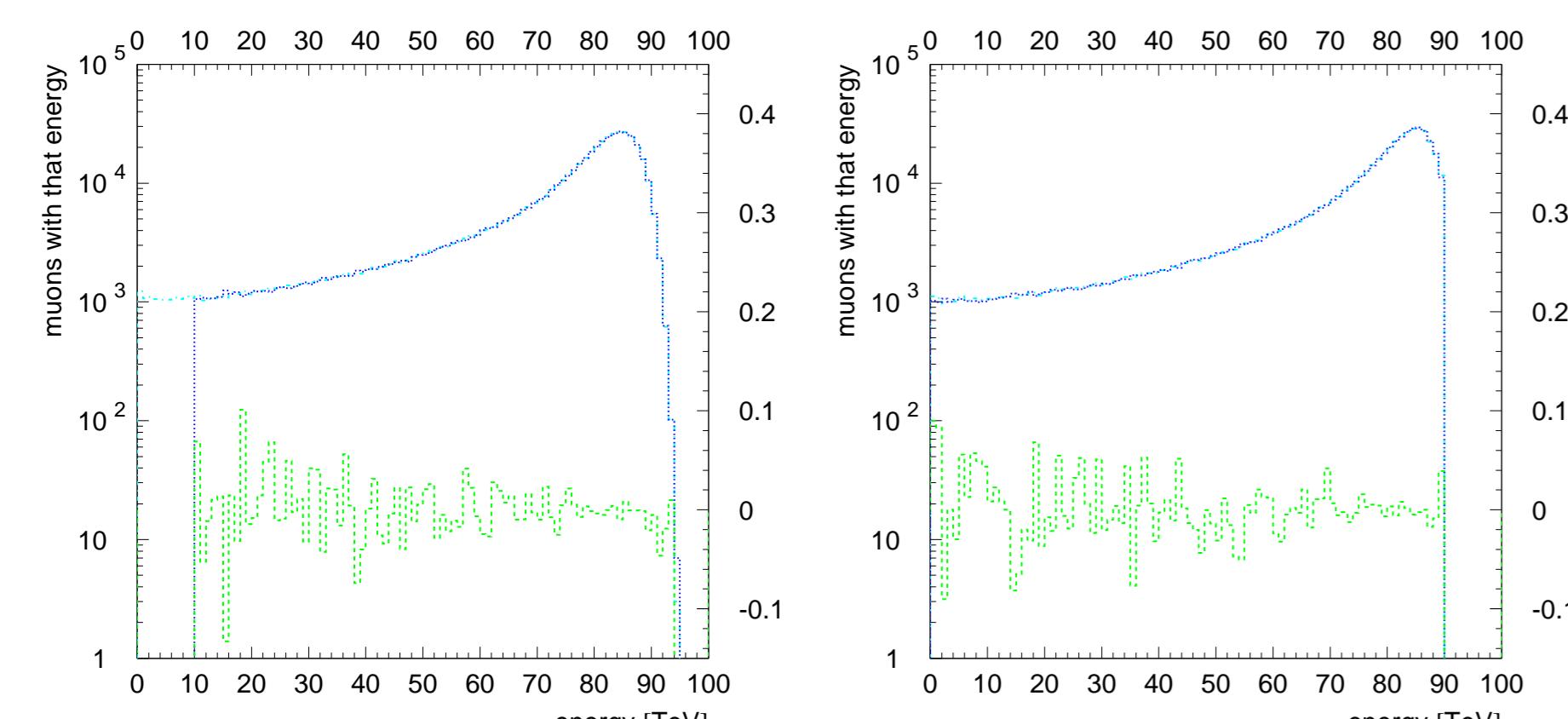


in Fréjus rock: solid line designates the value of the range evaluated with the second table (continuous and stochastic losses) and the broken line shows the range evaluated with the first table (continuous losses only).

Parametrization errors:



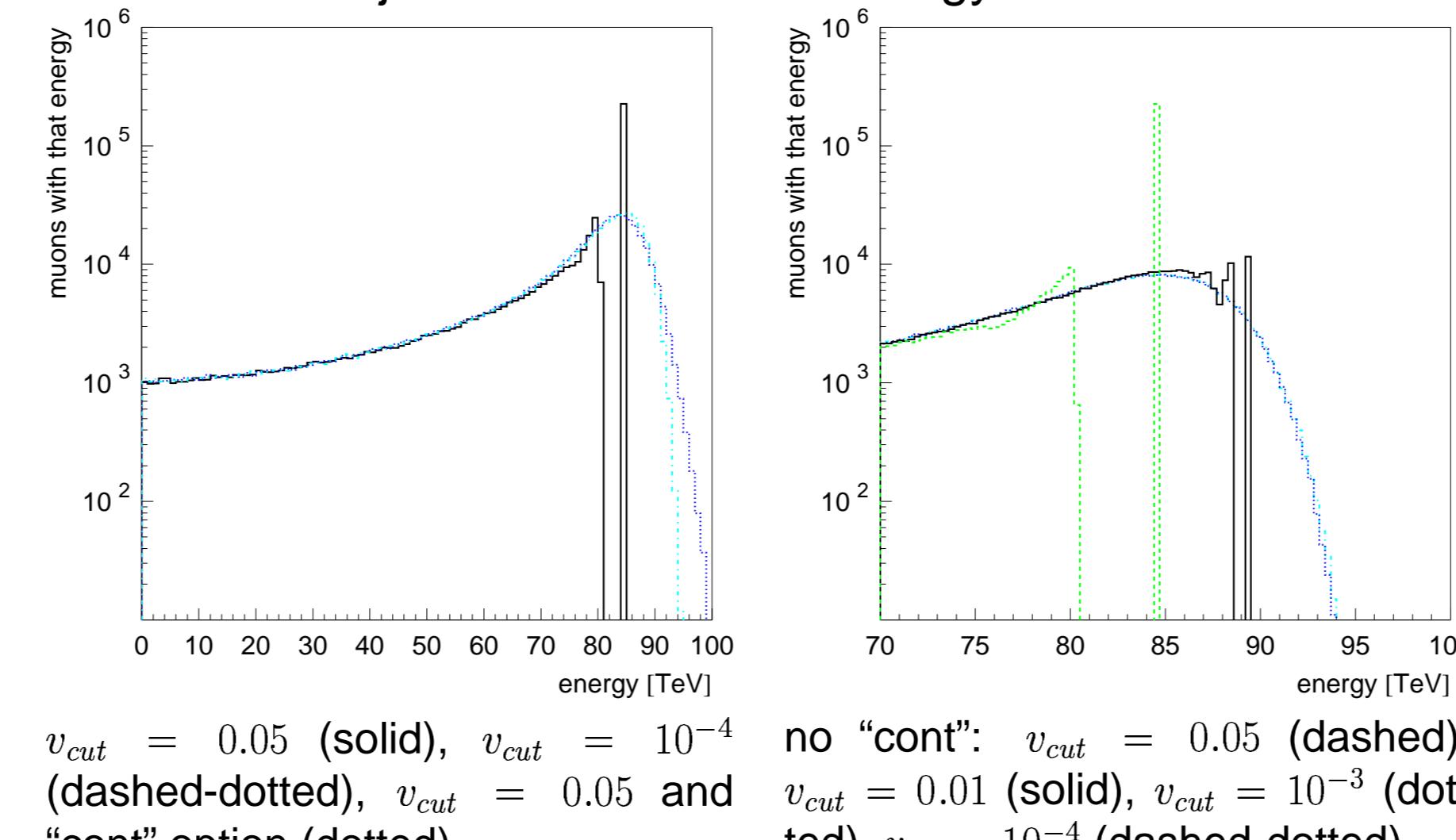
Final energy distribution of the muons that crossed 300 m of Fréjus Rock with initial energy 100 TeV:



Comparison of $e_{low} = m_\mu$ (dotted-dashed) with $e_{low}=10$ TeV (dotted). Also shown is the relative difference for $v_{cut} = 0.01$.

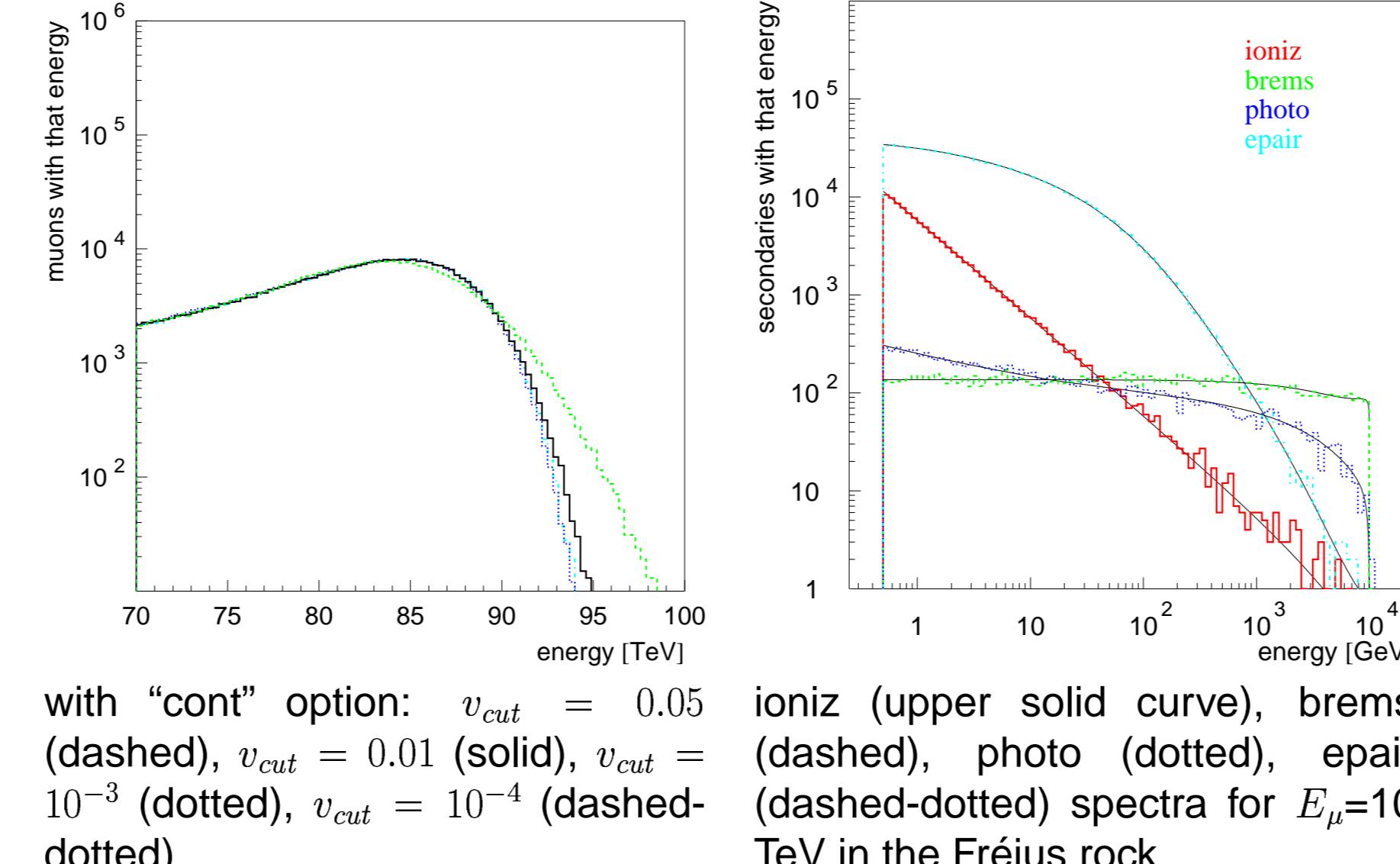
Comparison of parameterized (dashed-dotted) with exact (non-parametrized, dotted) versions for $v_{cut} = 0.01$.

Final energy distribution of the muons that crossed 300 m of Fréjus Rock with initial energy 100 TeV:



$v_{cut} = 0.05$ (solid), $v_{cut} = 10^{-4}$ (dashed-dotted), $v_{cut} = 0.05$ and "cont" option (dotted)

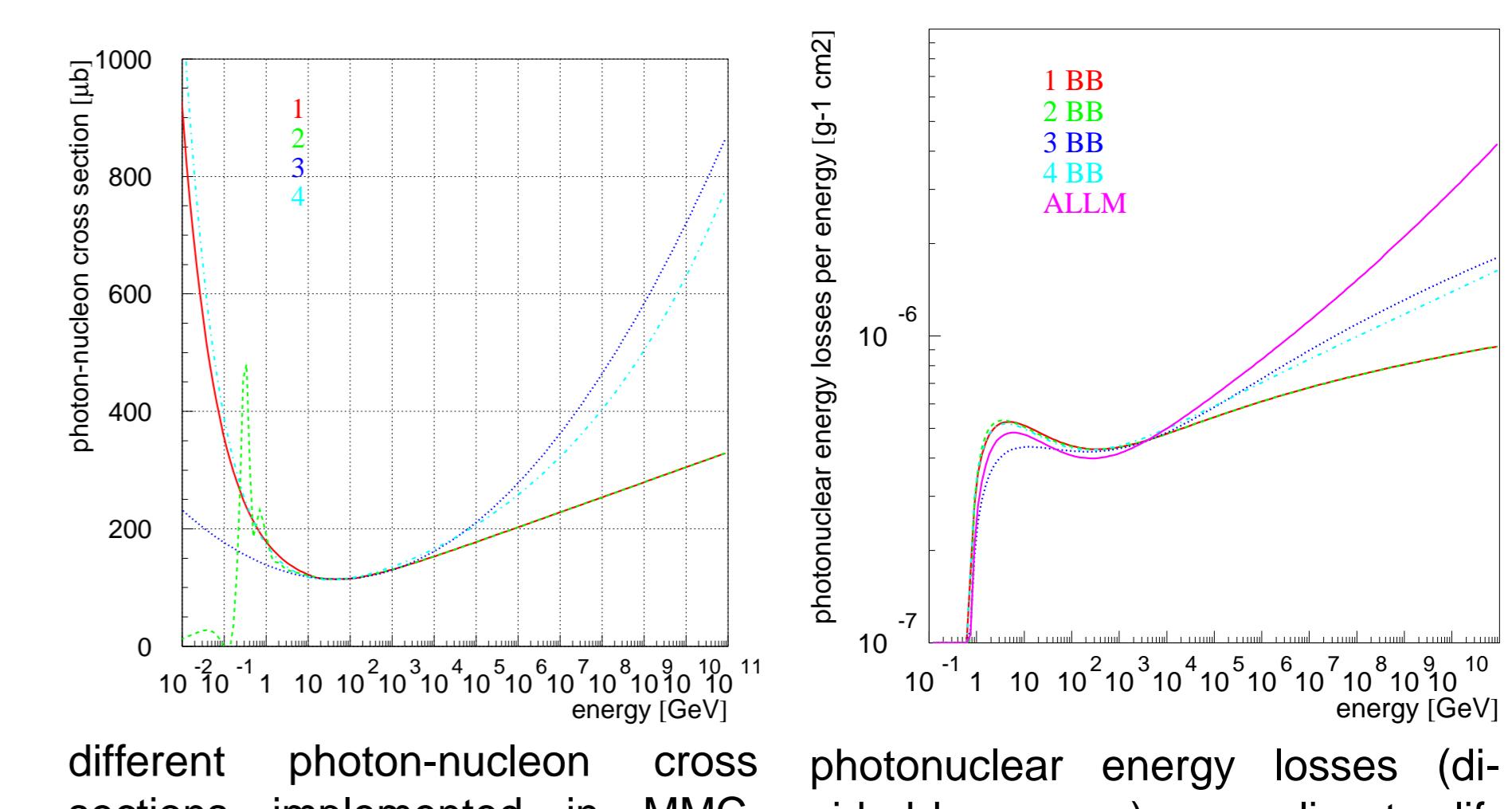
no "cont": $v_{cut} = 0.05$ (dashed), $v_{cut} = 0.01$ (solid), $v_{cut} = 10^{-3}$ (dotted), $v_{cut} = 10^{-4}$ (dashed-dotted)



with "cont" option: $v_{cut} = 0.05$ (dashed), $v_{cut} = 0.01$ (solid), $v_{cut} = 10^{-3}$ (dotted), $v_{cut} = 10^{-4}$ (dashed-dotted)

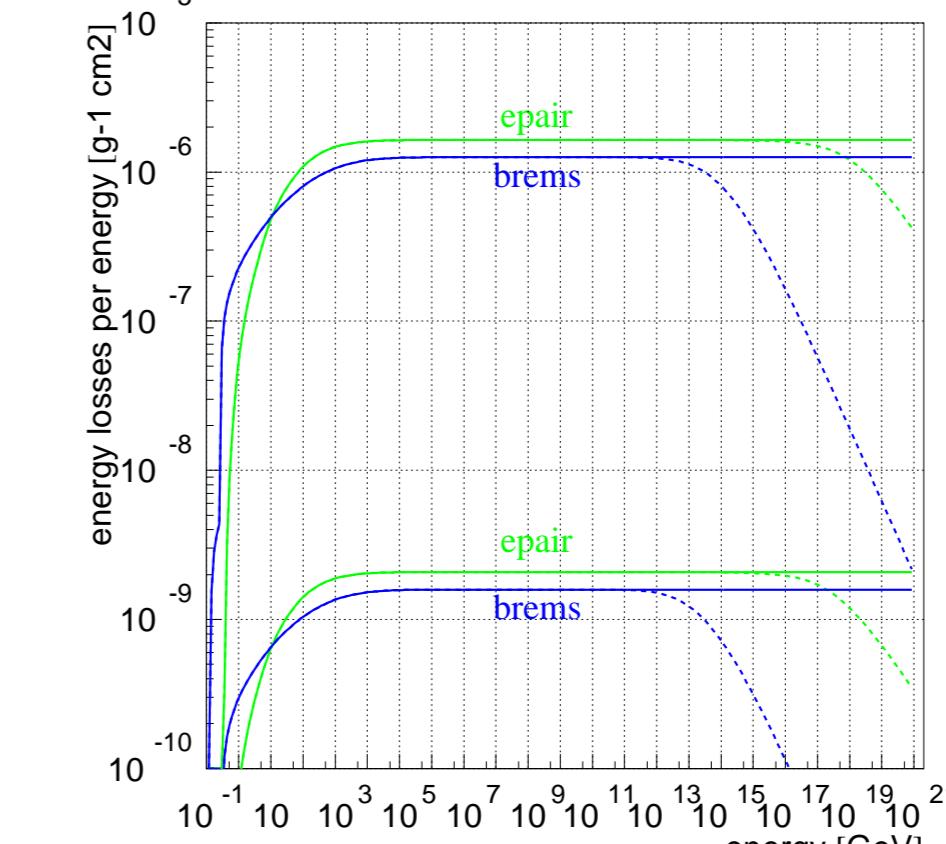
ioniz (upper solid curve), brems (dashed), photo (dotted), epair (dashed-dotted) spectra for $E_\mu=10$ TeV in the Fréjus rock

Photonuclear losses, LPM effect and Molière scattering

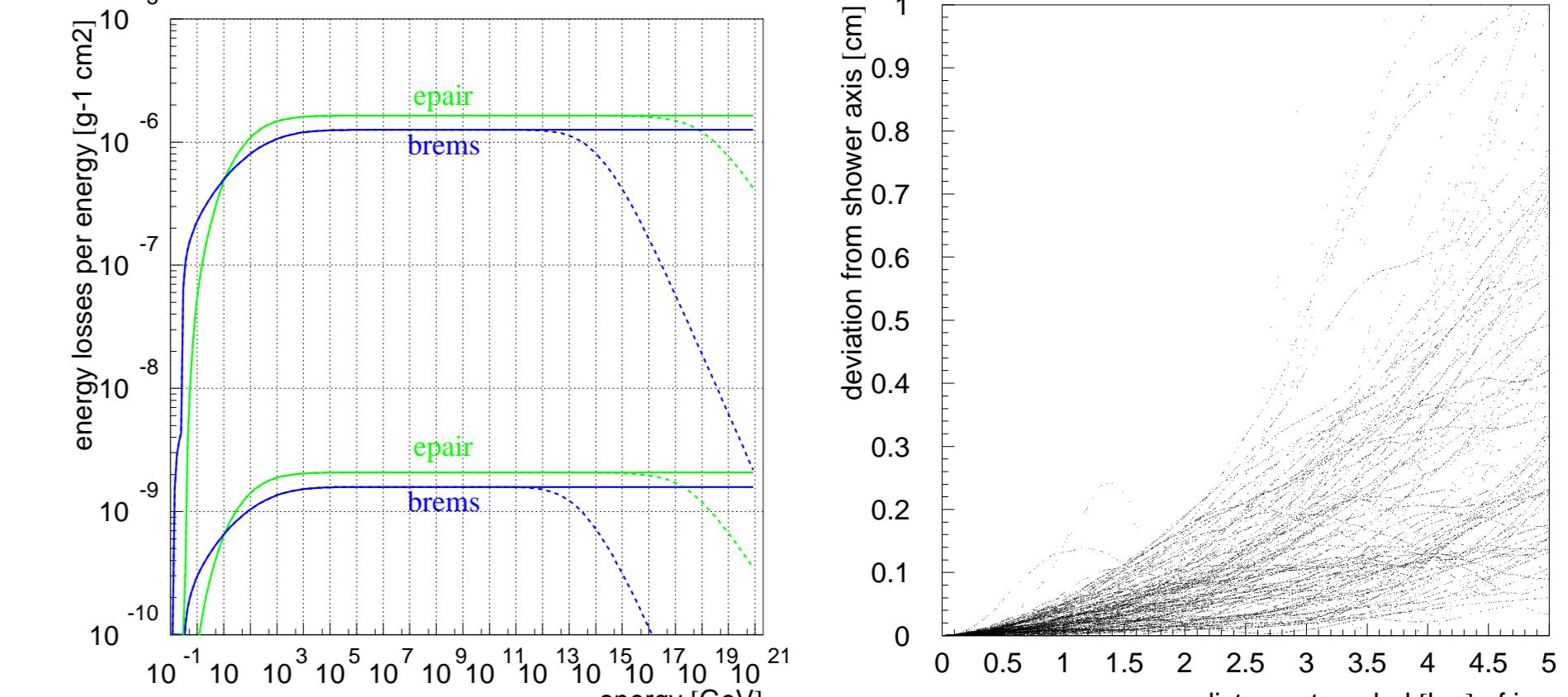


different photon-nucleon cross sections implemented in MMC, Bezrukov Bugaev parametrization

photonuclear energy losses (divided by energy), according to different formulae.



LPM effect in ice (higher plots) and Fréjus rock (lower plots, multiplied by 10⁻³)



Molière scattering of 100-10 TeV muons going straight down through ice