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TESTING ALTERNATIVE OSCILLATION SCENARIOS WITH ATMOSPHERIC NEUTRINOS USING AMANDA-II DATA FROM 2000 TO 2003

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The AMANDA-II neutrino telescope detects upward-going atmospheric muon neutrinos penetrating the Earth from the Northern Hemisphere via the Cherenkov light of neutrino-induced muons, allowing the reconstruction of the original neutrino direction. Due to the high energy threshold of about 50 GeV, the declination spectrum is minimally affected by mass-induced neutrino oscillations; however, alternative oscillation models predicting effects at high energy can be tested and constrained. Of particular interest are models that allow one to test Lorentz invariance and the equivalence principle. Using the AMANDA-II data from the years 2000 to 2003, a sample of 3401 candidate neutrino-induced events was selected. No indication for alternative oscillation effects was found. For maximal mixing angles, an upper limit is set on both the Lorentz violation parameter $\delta c/c$ and the equivalence level.

1 1. Introduction and detector description

² Cosmic ray particles entering the Earth's atmosphere generate a steady

- $_{3}$ flux of secondary particles, including muons and neutrinos. High energy
- ⁴ muons pass through the atmosphere and can penetrate several kilometers

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of ice and rock, while atmospheric neutrinos of energies only above roughly 5 40 TeV start to be absorbed in the Earth. Lower energy muon neutrinos penetrating the diameter of the Earth can oscillate into tau neutrinos. However, the oscillation maxima at $30 \,\mathrm{GeV}^1$ and below are beneath the 8 AMANDA-II threshold. Departures from conventional mass-induced oscillations could emerge at higher neutrino energies due to relativity-violating 10 effects (see below). Such mechanisms would distort the expected angular 11 distribution and energy spectrum of atmospheric neutrinos and could be 12 detectable by AMANDA-II. 13 The AMANDA-II neutrino telescope is embedded $1500 - 2000 \,\mathrm{m}$ deep 14 in the transparent and inert ice of the Antarctic ice sheet, close to the 15 geographic South Pole. AMANDA-II consists of 677 optical modules (OMs) 16

¹⁷ on 19 vertical strings, which are arranged in three approximately concentric ¹⁸ circles of 60 m, 120 m and 200 m diameter. Muons produced in ν_{μ} -nucleon ¹⁹ interactions can be directionally reconstructed by observing the Cherenkov ²⁰ radiation that propagates through the ice to the array of photosensors. To ²¹ ensure that the observed muon is due to a neutrino interaction, the Earth ²² is used as a filter against atmospheric muons, and only tracks from the

²³ Northern Hemisphere (declination $\delta > 0^{\circ}$) are selected.

Phenomenology of standard and alternative neutrino oscillations

It is commonly accepted that standard (mass-induced) $\nu_{\mu} \rightarrow \nu_{\tau}$ oscil-26 lations^a are responsible for the measured deficit of atmospheric muon 27 neutrinos¹. Atmospheric neutrino data can also be used to test non-28 standard oscillation mechanisms that lead to observable differences at 29 higher neutrino energies. Various new physics scenarios can result in neu-30 trino flavor mixing. Two of these scenarios, which can be described in a 31 mathematically analogous way, have been tested in this analysis. The un-32 derlying theories assume small deviations from the principles of the theory 33 of relativity and lead to measurable neutrino oscillations: 34

- In theories predicting violation of Lorentz invariance (VLI), a set of additional neutrino eigenstates with different maximal attainable
- velocities (MAV) c_n/c is introduced, violating special relativity².

^aIn the regime of atmospheric neutrino oscillations, it suffices to consider a two-flavor system of eigenstates (ν_{μ}, ν_{τ}).

38	• In theories predicting violation of the weak equivalence princi-
39	ple (VEP), gravitational neutrino eigenstates are introduced which
40	couple with distinct strengths γ_n to a gravitational potential $\phi,$ con-
41	flicting with the universal coupling assumed in general relativity ^{3,4} .

⁴² The main difference between these oscillation scenarios and standard oscil-⁴³ lations is the linear energy dependence of the oscillation frequency, shifting ⁴⁴ observable oscillation effects into the energy range of AMANDA-II. For the ⁴⁵ sake of simplicity, we will focus on the VLI scenario. As both theories ⁴⁶ are mathematically equivalent, the results can be transferred to the VEP ⁴⁷ case by simply exchanging the relativity-violating oscillation parameters ⁴⁸ $\delta c/c \rightarrow 2|\phi|\delta\gamma$ and mixing angles $\Theta_c \rightarrow \Theta_{\gamma}$.

⁴⁹ Combining standard and VLI oscillations, one obtains three systems ⁵⁰ of neutrino eigenstates (flavor, mass, and MAV eigenstates), resulting in ⁵¹ a total of 5 oscillation parameters: the mass-squared difference Δm^2 , two ⁵² mixing angles Θ_m and Θ_c , the VLI parameter $\delta c/c$, and a complex phase ⁵³ η . Fixing $\Delta m^2 = 2.3 \times 10^{-3} \,\mathrm{eV}^2$ and $\Theta_m = 45^\circ$, the survival probability ⁵⁴ may then be written as:

$$P(\nu_{\mu} \to \nu_{\mu}) = 1 - \sin^2 2\Theta \, \sin^2 \left(\Omega \, L\right) \tag{1}$$

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 $2\Theta = \arctan\left(s/t\right) \qquad \Omega = \sqrt{s^2 + t^2} \tag{2}$

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$s = 2.92 \times 10^{-3} \left 1/E_{\nu} + 8.70 \times 10^{20} \ \delta c/c \ \sin 2\Theta_c \ E_{\nu} \ \mathrm{e}^{i\eta} \right $,
$t = 2.54 \times 10^{18} \ \delta c/c \ \cos 2\Theta_c \ E_{\nu}$.	(3)

Here the the muon neutrino path length L is expressed in km and the neutrino energy E_{ν} in GeV. For the given set of parameters, one can observe a significant effect within the analyzed energy range (100 GeV - 10 TeV)and declination range ($\delta \geq 20^{\circ}$), for certain values of Θ_c and $\delta c/c$.

The VLI effects can also be viewed in terms of the Standard Model 66 Extension, an effective quantum field theory which adds all coordinate-67 independent renormalizable Lorentz-violating operators to the Standard 68 Model^{5,6}. In particular, the above oscillations correspond to non-zero and 69 differing eigenvalues of c_L , the dimensionless coefficient of the Lorentz-70 violating Lagrangian term $\frac{1}{2}i(c_L)_{\mu\nu ab}\overline{L}_a\gamma^{\mu}\overline{D}^{\nu}L_b$ ⁷. Additionally, in this 71 scenario only rotationally invariant (time) components of c_L are non-zero 72 (so-called "fried chicken" models), and we consider only a two-flavor system. 73 With these restrictions, $\delta c/c$ in eq. 3 is equivalent to the difference in 74 eigenvalues $(c_a - c_b)$ of the 2 × 2 matrix c_L^{TT} .

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⁷⁶ 3. Data selection and analysis method

The data analyzed in this analysis are selected from 7.9×10^9 events recorded 77 from 2000 to 2003, representing a total livetime of 807.2 days. Quality se-78 lection criteria based on various track reconstruction quantities are used 79 to reduce contamination of the sample by mis-reconstructed atmospheric 80 muons to a level of approximately 4% (see ref. 8 for more details). The re-81 sulting number of selected neutrino candidate events is 3401. Reconstructed 82 declination and the number of OMs triggered in an event $(N_{\rm ch}, \text{ an energy})$ 83 estimator) are used as observables to search for any deviations in the atmo-84 spheric neutrino spectrum that are dependent on flight length or neutrino 85 energy. A full simulation chain, including neutrino absorption in the Earth, 86 neutral current regeneration, muon propagation, and detector response is 87 used to simulate the response of AMANDA-II to atmospheric neutrinos^{9,10}. 88 The expected atmospheric muon neutrino flux before oscillations is taken 89 from Lipari¹¹. 90

The analysis method uses a χ^2 -test to compare the declination and $N_{\rm ch}$ 91 distributions of data with Monte Carlo simulations including VLI oscillation 92 effects. The systematic uncertainties affecting the Monte Carlo prediction 93 are integrated into the χ^2 expression: uncertainties in the normalization 94 due to detector response and theoretical models of atmospheric neutrino 95 flux (30%); uncertainty due to the relative production rate between kaons 96 and pions (6%); and uncertainty in the sensitivity of the OMs (11.5%). 97 The optimal binning of the χ^2 expression is determined using toy Monte 95 Carlo studies, and the exclusion regions for alternative oscillation effects 99 are obtained by scanning through the oscillation parameter space. 100

101 4. Results and Outlook

The analysis of the final atmospheric neutrino sample finds no evidence for 102 alternative oscillations, and a preliminary upper limit on the VLI parameter 103 $\delta c/c$ is set of 5.3×10^{-27} at the 90% confidence level, for nearly maximal 104 mixing angles $\Theta_c \approx \pm \pi/4$. The dependence on the unconstrained phase η 105 is found to be small (see figure 1); the most conservative limit is obtained 106 for $\cos \eta = 0$. The limit can also be interpreted in the context of VEP 107 theories, leading to an upper limit of $2|\phi|\delta\gamma \leq 5.3 \times 10^{-27}$. This result is 108 competitive with the limits obtained using data from Super-Kamiokande¹² 109 and MACRO¹³. However, AMANDA is not sensitive to small mixing angles 110 due to the systematic errors and its higher energy threshold. 111

A likelihood analysis of the 2000-2006 AMANDA-II data sample is

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Figure 1. Shown are preliminary exclusion regions for VLI (VEP) oscillation effects, left for $\cos \eta = 0$, right for $\cos \eta = 1$.

¹¹³ in progress, with improved systematic error estimation and increased ¹¹⁴ sensitivity¹⁴. This analysis will also extend the technique to search for ¹¹⁵ evidence of quantum decoherence resulting from interaction of neutrinos ¹¹⁶ with the background space-time foam¹⁵. The next-generation IceCube de-¹¹⁷ tector, when completed in 2010, will be able to extend the sensitivity to ¹¹⁸ VLI effects by about one order of magnitude¹⁶.

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